

Comments on Self-Focusing Models of Stationary and Non-Stationary Optical Beams

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Introduction

- **Self-focusing (SF)** of optical beams as one of fundamental phenomena in nonlinear optics **attracts a great interest** of researches in quantum electronics and laser physics
- **Reasons** for this:
 - **A strong influence** of SF on interaction processes of high-power laser radiation with optical media
 - **Potentials** for **possible applications** (remote atmospheric sounding, electric discharge control, etc)
- **Investigation of SF mechanisms** in optical media and elucidation of adequate SF models – important directions in laser physics and nonlinear optics
- A huge number of **experimental and theoretical works** done in these directions during almost 50 years **resulted in understanding major features of SF** phenomenon, various SF models proposed and investigated theoretically and experimentally

Introduction (continued)

- However, **some questions** regarding these features and models (especially in ultra-short time domain) remain **unclear** and require further studies
- In this context I present here **comments** on the **SF models** published to date, and **discuss** some **prospects** for future studies
- **Plan of the report:**
 - A brief historical **review** of **SF studies**: major stage, theoretical models, experiments
 - A brief **description** of **multi-focus structure** and **moving nonlinear foci models** as the most adequate models of SF
 - A brief **review** and **comments** on **SF of fs-pulses**
 - A **discussion** of **prospects** for further studies

History of SF: major stages

- **Prediction of SF phenomenon**, qualitative analysis, proposal for wave-guiding model (Askaryan, 1962)
- **First observations of SF** (Hercher, 1964; Pilipetski and Rustamov, 1965)
- **First comprehensive theoretical studies of SF**, proposal for **self-trapping model** (Chiao, Garmire and Townes; Talanov, 1964)
- **Further theoretical studies:**
 - Determination of main **characteristics of SF**: critical power, SF length, self-phase modulation, spatial instability (Bespalov, Talanov, Akhmanov, Khokhlov, Sukhorukov, Kelly, Marburger, 1964-1972)
 - **Proposal for alternative SF models: multi-focus structure** and (Dishko, Lugovoi, and Prokhorov, 1967), **moving nonlinear foci** (Lugovoi and Prokhorov, 1969)
- Purposeful **experimental studies of SF** confirming **adequacy of MFS-MNLF models** (Loy and Shen, 1969; Korobin, et al, 1970; Manenkov, et al, 1970)

History of SF (continued)

—A new stage in SF studies, fs-pulses:

- **First observation** of SF of fs-pulses (**filamentation**) in **air** (Mourou, et al, 1995)
- **Subsequent studies:**
 - Elucidation of **main characteristics** of **filamentation**: length, diameter, and energy of filaments, transformation of spectrum-supercontinuum and conical emission (Mysyrowicz, et al, 1996; Chin, et al, 1997)
 - Theoretical studies, **proposals** for **filamentation models**: self-guiding, self-trapping, moving foci and refocusing, spatial replenishment (Braun, et al; Mysyrowicz, et al; Kandidov, et al; Molony, et al, 1997)
 - Continuing studies of fs-pulses filamentation in air, other gases, and condensed media (many researchers in many countries, 1997- up to now)

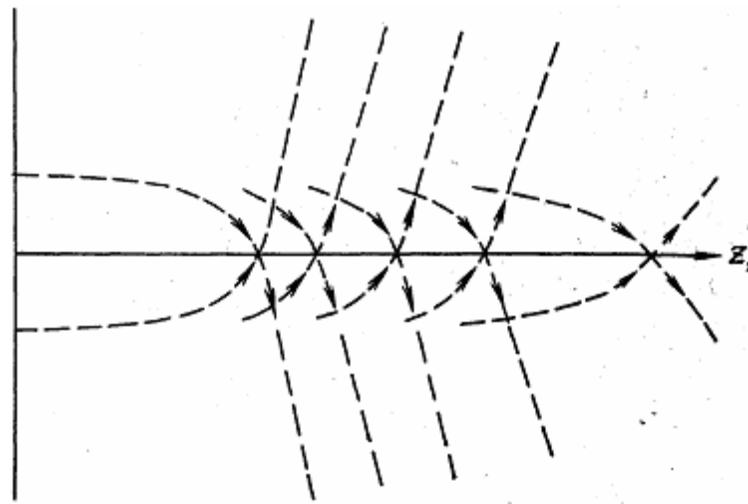
MFS-MNLF models

- Analysis of SF studies shows: **main results** can be **explained** most **adequately** on a basis of **MFS-MNLF models**
- Below I describe briefly a **theory of these models** and **some experimental results** confirming their adequacy
- **MFS model** has been **formulated** (Dyshko, Lugovoi, and Prokhorov, 1967) on a base of numerical investigation of a light beam propagation in **Kerr-type nonlinear medium** described by **nonlinear wave equation**:

$$\frac{\partial^2 E}{\partial r^2} + \frac{1}{r} \frac{\partial E}{\partial r} + 2ik \frac{\partial E}{\partial z} + n_2 k^2 |E|^2 E = 0 \quad (1)$$

- Established: a **series of nonlinear foci** is formed on a beam propagation axis, Z:

Fig.1



MFS-MNLF models (continued 1)

- **Location of nonlinear foci:**

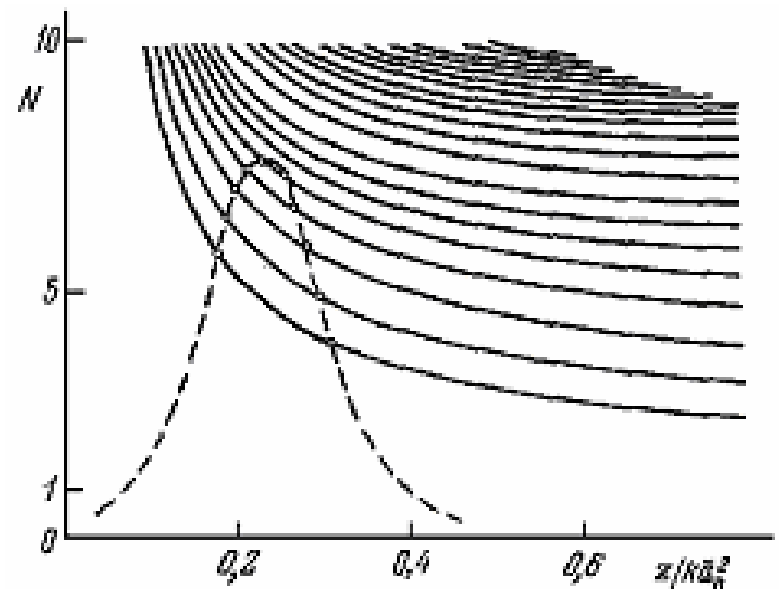
$$\xi_{\Phi_m} = \frac{\chi_m}{N_m} \frac{k\bar{a}_0^2}{\sqrt{P_0/P_{cr}^{(m)} - 1}} \quad (2)$$

P_0 and a_0 - incident beam power and radius, respectively, $P_{cr}^{(m)}$ is critical power for m-th focus $P_{cr}^{(m)} \cong mP_{cr}^{(1)}$, $P_{cr}^{(1)} = cn_0 N_1^2 / 8n_2 k^2$,

—For a **stationary light beam** (time independent power $P=P_0$) the **location of nonlinear foci is fixed** on Z-axis whereas for a **non-stationary beam** ($P=P(t)$) **foci move** along Z-axis

—Fig.2 illustrates **formation of moving multi-focus structure** in a fast-response Kerr-medium. Here **solid curves** - location of nonlinear foci at various pulse power, **dashed curves** - light pulse shape

Fig. 2

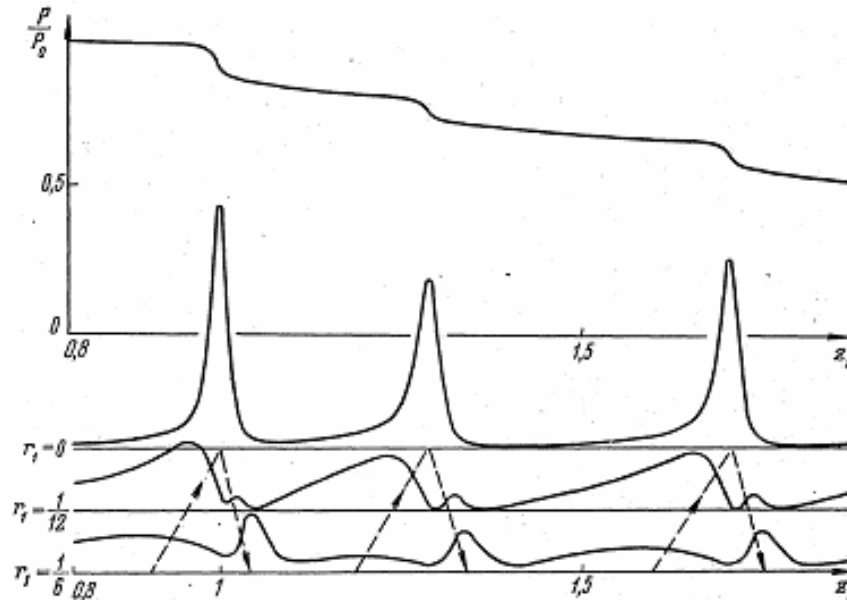


MFS-MNLF models (continued 2)

— Theoretically established:

- Only a portion of an initial beam power, $\cong P_{cr}^{(1)}$, flows through each nonlinear foci, a rest power of the beam is absorbed and diffracts in the foci at rather large angles
- Longitudinal and transverse sizes of the foci, as well as their location on Z-axis, depend on physical processes in the media (nonlinear absorption, ionization, etc)
- However, multi-focus structure is, in general, remained, i.e. is stable to perturbations of nonlinear properties of the media (Fig. 3 illustrates this fact for a medium with 3-photon absorption)

Fig.3



MFS-MNLF models (continued 3)

- MFS model is stable also to variation to spatial beam profile (super-Gaussian beams were theoretically investigated, Danileiko, et al, 1981)
- Multiple filamentation, observed in case of inhomogeneous beams, is naturally explained in a frame of MFS-MNLF models: separate MFS are formed if intensity spikes of the beam exceed P_{cr} .

MFS-MNLF models: experiment

- At **earlier stages** of SF studies (**1962-1970**) observed **filaments** were interpreted in a frame of **self-waveguide-self-trapping** concept
- In **1969-1970** **purposeful experimental studies** have been performed to prove what models, self-waveguide or moving nonlinear foci, are adequate:
 - Observation of spatial-temporal **evolution** of a **laser beam** at the **exit** and **inside** of a nonlinear optical liquid cell (Loy and Shen, 1969)
 - Observation of spatial-temporal **evolution** of the **laser beam** in a nonlinear optical liquid cell applying **high-speed temporal resolution technique** (Korobkin, et al, 1970)
 - Observation of spatial structure of a laser beam in a nonlinear medium (glass) at **variation** of **temporal laser pulse shape** (Lipatov, Manenkov, and Prokhorov, 1970)

MFS-MNLF models: experiment (continued 1)

- All these experiments have confirmed adequacy of the moving nonlinear foci model
- Here we present the results of our experiment (Lipatov, Manenkov, Prokhorov, 1970) in which self-focusing of ruby laser beam in glass has been investigated at different temporal pulse shapes: bell-like, triangular, and rectangular (flat-top)
 - In cases of bell-like and triangular pulse shapes filamentary damage in the glass has been observed, whereas in the case of rectangular pulse shape the point-like damage has been observed.
 - Fig.4 illustrates the latter case:
laser-induced light scattering patterns at various incident powers:
(a) $P_i = P_d$, (b) $P_i = 3P_d$, (c) $P_i = 6P_d$

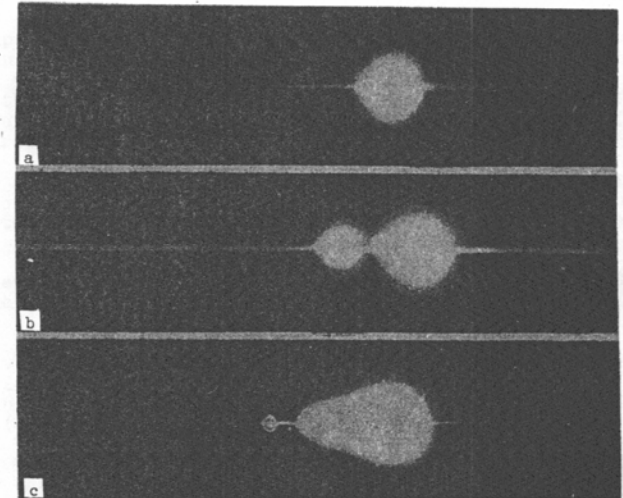


Fig. 4

MFS-MNLF models: experiment (continued 2)

- Interpretation of the results:
 - **Filamentary damage** is caused by **non-stationary SF** (moving nonlinear foci)
 - **Point-like damage** is the result of **stationary SF** (standing nonlinear foci)

SF of fs-pulses

- A huge number of experimental and theoretical works has been carried out on SF of fs-pulses during last 15 years
- Main results of experimental works:

In air:

- Strongly pronounced beam filamentation is observed at $P_o \geq P_{cr}$ ($\approx 3\text{GW}$): very long ($l_f \geq 10\text{ m}$) and thin ($d_f \cong 100\ \mu\text{m}$) laser light and plasma filaments
- Light energy in the filament $E_f \cong 6\%P_o$
- Multiple filaments are observed at $P_o \geq 10P_{cr}$
- A strong transformation of spectrum is observed at $P_o > P_{cr}$:
 - large broadening: super-continuum (230 nm-4 μm)
 - Specific conical emission

In condensed media:

- P_{cr} is an order of magnitude less than in gases,
- $l_f \cong$ several cm,
- $d_f \cong 2\ \mu\text{m}$

SF of fs-pulses(continued 1)

- Theoretical results :
- Theoretical studies of fs-pulse filamentation (especially in air) have been performed by several groups in Russia, US, France, Canada, and other countries
- All these studies were based (with some variations) on **numerical solutions of nonlinear wave equation** (nonlinear Schrodinger equation) taking into account ordinary light diffraction and several nonlinear effects: **Kerr-nonlinearity** of medium refractive index, **multi-photon absorption (MPA)**, and **plasma generation** due to **MPA**

$$2i \frac{\partial E}{\partial z} + \frac{1}{k_0} \Delta_{\perp} E - k'' \frac{\partial^2 E}{\partial t^2} + k_0 n_2 (|E|^2 + \tau_K^{-1} \int_{-\infty}^1 e^{-(t-t')/\tau_K} |E(t')|^2 dt') E - k_0 \frac{\omega_{pe}^2(\rho)}{\omega_0^2} E + i\beta^{(K)} |E|^{2K-2} E = 0 \quad (3)$$

$$\frac{\partial \rho}{\partial \tau} = \frac{\beta^{(K)}}{K\hbar\omega_0} |E|^{2K} \left(1 - \frac{\rho}{\rho_{at}} \right) \quad (4)$$

SF of fs-pulses (continued 2)

- Computer **simulation results** have been **interpreted** on a base of different models:
 - **Self-channeling** (Braun, et al, 1997)
 - **Dynamic spatial replenishment** (Molonely, et al, 1997)
 - **Moving foci and refocusing** (Kandidov, et al, 1997)
 - **Slice-by-slice self-focusing** (Chin, et al, 1997)
- An analysis (Lugovoi and Manenkov, 2005) of these computer simulations has shown:
 - **some shortcomings**: incomplete description of group velocity dispersion- not taking into account a plasma component, inadequate interpretation of delayed response term in **Kerr - nonlinearity** of refractive index
 - **SF models suggested** coincide(in essence) with Lugovoi-Prokhorov MFS-MNLF models although **terminology** used **is different**

SF of fs-pulses (continued 3)

— Indeed, for example:

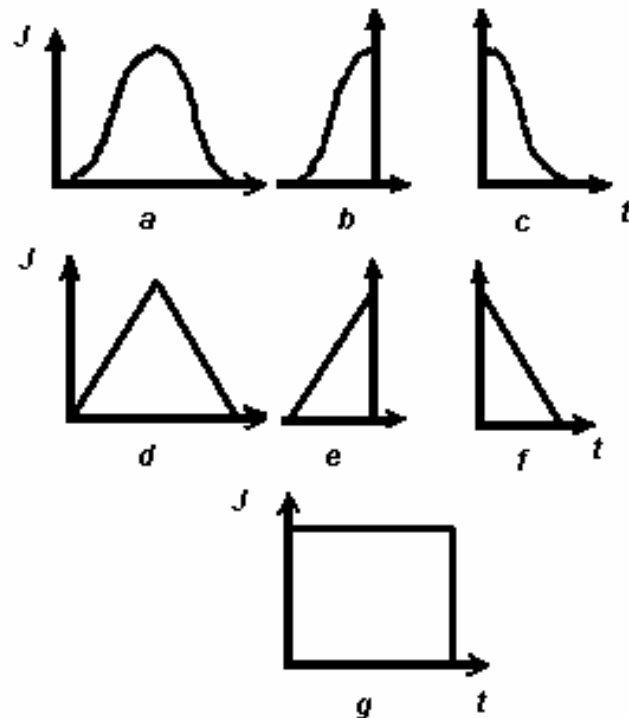
- “**self-guided light strings**” (Moloney, et al), in fact, correspond to “moving nonlinear foci” in Lugovoi-Prokhorov model
- “**nonlinear robust mode**” correspond to a portion of the beam flowing through the nonlinear focus in Lugovoi-Prokhorov model
- “**dynamic spatial replenishment**” corresponds to the process of multi-focus structure formation in Lugovoi-Prokhorov model
- Similarly, “**slice-by-slice self-focusing**” (Chin, et al) corresponds to multi-foci structure formation in Lugovoi-Prokhorov model

SF of fs-pulses(continued 4)

- Therefore, available **experimental** and **theoretical results** on **SF of fs-pulses** are **in a qualitative agreement** with **Lugovoi-Prokhorov MFS-MNLF models**
- At the same time, taking into account a complexity of processes involved into SF of fs-pulses (strong influence of plasma formation, etc), **further studies**, both experimental and theoretical, **are required** using new approaches

SF of fs-pulses: prospects

- Use of laser pulses with **variable temporal shapes** (similarly to that used in studies of SF of ns-pulses) seems to be **promising approach** to study and control SF processes
- Indeed, an analysis (Manenkov, 2008) has shown:



- If **bell-like** (e.g. Gaussian) or **triangular pulse shapes** are used (symmetric -Fig. 5a,d and asymmetric -Fig. 5b,c,e,f) **different** filamentary spatial **structures** of the beam intensity should be observed because of **different effects** of laser-induced **plasma on rise and fall times** of the pulse propagating in a medium

Fig. 5

SF of fs-pulses: prospects (continued 1)

- All **bell-like** and **triangular pulse** shapes correspond to **non-stationary SF**
- Quite **different beam structures** should be observed if **rectangular (flat-top) laser pulse shape** is used: a spatially standing discrete nonlinear foci structure (**no filamentation!**) corresponding **stationary SF**
- **Spectral properties** of optical radiation propagating in nonlinear media have to be also **significantly different** at various pulse shapes because of different character of spectral broadening due to **self-phase modulation**:

$$\Delta\omega = -\frac{d\varphi_{nl}}{dt} = -\frac{d}{dt}\left(\frac{\omega}{c}n_2I(t) \cdot L\right) \quad (5)$$

φ_{nl} – nonlinear phase

ω – central frequency of laser radiation

n_{nl} – nonlinear refractive index

$I(t)$ – intensity

L – interaction length

SF of fs-pulses: prospects (continued 2)

- In cases of **bell-like and triangular pulses** **different spectral broadening** should be observed depending on $I(t)$ shape
- In particular, if **asymmetric pulse shapes** are used **blue-shifted** or **red-shifted** spectral **broadening** should be observed
- In case of **rectangular pulse** spectral **broadening** due to self-phase modulation **should be absent** ($d\varphi_{nl}/dt=0$)

Conclusions

- **SF phenomenon** of optical beams has been **studied** to date **rather well** , especially in ns-ps pulse duration domain
- **Main features** of the phenomenon observed in ns-ps range are **explained** most **adequately** on the basis of **Lugovoi-Prokhorov MFS-MNLF model**
- This model **can**, we believe, also **explain** features of **SF in shorter (fs) pulse duration range**
- However, for more complete understanding of SF in fs-range **further** experimental and theoretical **studies** are **required**
- To our view, **a perspective approach** for such studies is **to use laser pulses with variable temporal shapes**